Study of Microwave Simulation Technology

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Resonant Frequency and Quality Factor Analysis with a Cylindrical FDTD Resonant frequency and unloaded (conductor, dielectric) Q values can be calculated for cylindrical cavities and dielectric resonators. The computation time can be reduced with an FDTD/GPOF combination.

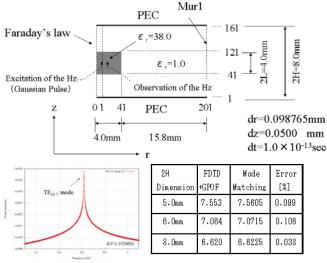
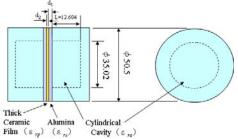


Table 1. Comparison of calculated results [GHz] and error rates

Analysis of a Cylindrical Cavity Resonator Using the Mode-matching Method The resonant frequency of a divided TE_{011} cylindrical cavity resonator can be calculated by inserting sample materials.



Resonant frequency is determined to satisfy detZ = 0 by applying the mode-matching method for the surround of a sample material with perfect electric conductor (PEC) properties, where Z is the matrix, and

$$Z = \begin{bmatrix} Q & -R \\ S & -P \end{bmatrix}$$
(1)

In this matrix element, m and n terminate as M and N as finite values. Moreover, Q, R, S and P are as follows:

$$Q_{mn} = U_n \frac{a \cdot ku_n}{ks_m^2 - ku_n^2} J_0(ku_n a) J_1(ks_m a) \cdot sin(\beta u_n \cdot L)$$
(2)

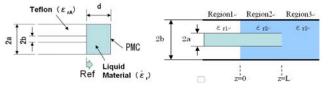
$$R_{mm} = V_m \frac{b^2}{2} J_0^2(ks_m b) \cdot \left[B'_m \cos\left(\beta s_m \cdot \frac{d_1}{2}\right) + C'_m \sin\left(\beta s_m \cdot \frac{d_1}{2}\right) \right]$$
(3)

$$S_{nn} = U_n \beta u_n \frac{a^2}{2} J_0^2(ku_n a) \cdot \cos(\beta u_n L)$$
(4)

$$P_{nm} = V_m \beta s_m \frac{a \cdot ku_n}{ks_m^2 - ku_n^2} J_1(ks_m a) J_0(ku_n a) \cdot$$

$$\left[B'_{m}\sin\left(\beta s_{m}\cdot\frac{d_{1}}{2}\right)-C'_{m}\cos\left(\beta s_{m}\cdot\frac{d_{1}}{2}\right)\right]$$
(5)

Discontinuity Analysis at the Boundary Plane from the Coaxial to the Cylindrical Waveguide The input impedance of a discontinuity structure from the coaxial to the circular waveguide can be calculated if each region's dielectric constant is different.



Using the mode-matching method, the input admittance can be calculated using Equation (6):

$$\dot{\mathbf{y}}_{m} = j \cdot \frac{2k_{0}a \cdot \dot{\boldsymbol{\varepsilon}}_{r}}{\sqrt{\boldsymbol{\varepsilon}_{rA}} \cdot \ln\left(\frac{a}{b}\right)} \left(\mathbf{y}_{0} - \sum_{q=1}^{\infty} \boldsymbol{y}_{q} \cdot \boldsymbol{x}_{q} \right)$$
(6)

where

$$y_{0} = \sum_{i=1}^{\infty} \frac{tanh[(\lambda_{i}^{2}a^{2} - k^{2}a^{2})^{1/2}d/a]}{(\lambda_{i}^{2}a^{2} - k^{2}a^{2})^{1/2}\lambda_{i}^{2}a^{2}} \cdot \frac{J_{0}^{2}(\lambda_{i} \cdot b)}{J_{1}^{2}(\lambda_{i} \cdot a)}$$
(7)

and
$$x_q$$
 is calculated by

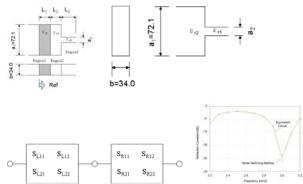
 $\boldsymbol{x}_q = \boldsymbol{A}_{qn}^{-1} \cdot \boldsymbol{y}_q \cdot$

In their matrix element, A_{qn} is

$$A_{qn} = \sum_{i=1}^{\infty} \frac{\lambda_i^2 a^2 \cdot tanh[\lambda_i^2 a^2 - k^2 a^2]^{1/2} d/a]}{(\lambda_i^2 a^2 - k^2 a^2)^{1/2} (\lambda_i^2 a^2 - \xi_q^2 a^2) \cdot (\lambda_i^2 a^2 - \xi_n^2 a^2)} \cdot \frac{J_0^2(\lambda_i \cdot b)}{J_1^2(\lambda_i \cdot a)} + \frac{1}{\dot{\varepsilon}_r} \cdot \frac{\delta_{qn}}{4(\xi_n^2 a^2 - k_0^2 a^2 \cdot \varepsilon_{rA})^{1/2}} \cdot \left[\frac{a^2}{b^2} \cdot \frac{Z_1^2(\xi_n \cdot a)}{Z_1^2(\xi_n \cdot b)} - 1\right]$$
(9)

(8)

Rectangular Waveguide Discontinuity Analysis Using the Generalized Scattering Matrix The S-parameter of the rectangular waveguide discontinuity can be calculated.



The mode-matching method can be applied to this analytical model. Accordingly, the S-parameter for a single h-plane discontinuity of the rectangular waveguide as shown in Fig. 4 can first be calculated as follows:

$$S_{11} = [L_E L_H + I]^{-1} [L_E L_H - I] \dots (10)$$

$$S_{12} = 2[L_E L_H + I]^{-1} L_E \dots (11)$$

$$S_{21} = L_H \{I - [L_E L_H + I]^{-1} [L_E L_H - I]\} \dots (12)$$

$$S_{22} = I - 2L_H [L_E L_H + I]^{-1} L_E \dots (13)$$